

Suspension Modelling in the Lattice Boltzmann Method

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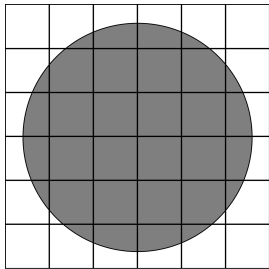
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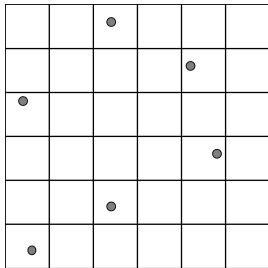
- Introduction
- Non resolved approach
 - Coupling approaches
- Resolved approach
 - Volume approximation techniques
 - Partially saturated method (PSM)
 - Homogeneous lattice Boltzmann method (HLBM)
 - Benchmark comparison
- Conclusion

Resolved and Non-resolved Approaches



Resolved approach

- Particle diameters d_p are greater than the cell length Δx
- Particles are simulated by different approaches that recover the shape



Non resolved approach

- Particle diameters d_p are smaller than the cell length Δx
- Particles are modelled by adding forces

- Forced lattice Boltzmann equation

$$f_i(\mathbf{x} + \mathbf{c}_i \Delta t, t + \Delta t) = f_i(\mathbf{x}, t) + \Omega_i(\mathbf{x}, t) + \Delta t F_i(\mathbf{x}, t)$$

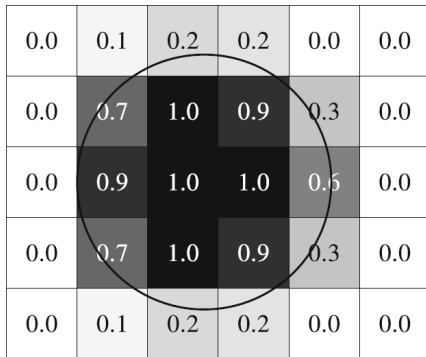
- Newton's law of motion

$$m_s \frac{d\mathbf{u}_s}{dt} = m_s \frac{d^2 \mathbf{x}_s}{dt^2} = F(\mathbf{x}, t)$$

- Volume fraction of solid phase:
 - **Dilute:** Low volume fraction, Impact of particles on the fluid flow can be neglected
 - **Dense:** High volume fraction, the fluid flow is affected by the presence of the particle
- One-way-coupling
 - Hydrodynamic forces acting on particles
- Two-way-coupling
 - + Particles forces acting on the fluid
- Four-way-coupling
 - + Particle-particle interaction (collision model)
 - + Particle wall interaction

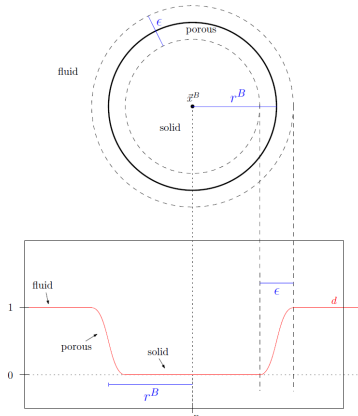
- Solid volume fraction $\varepsilon(\mathbf{x}, t)$

$$\varepsilon(\mathbf{x}, t) = \frac{V_s(\mathbf{x}, t)}{V_{total}(\mathbf{x}, t)}$$



- Estimation of the solid volume fraction $\varepsilon(\mathbf{x}, t)$ is difficult
- Compromise of computing time and accuracy
- Established methods:
 - Exact closed form solution
 - Polyhedral approximation
 - **Smoothed profile method**
 - **Supersampling**

$$\varepsilon(\mathbf{x}, t) = \begin{cases} 0 & \|\mathbf{x} - \mathbf{x}^B(t)\|_2 - r^B + \epsilon/2 \geq \epsilon \\ \sin^2 \left(\frac{\pi(\|\mathbf{x} - \mathbf{x}^B(t)\|_2 - r^B + \epsilon/2)}{2\epsilon} \right) & \|\mathbf{x} - \mathbf{x}^B(t)\|_2 - r^B + \epsilon/2 \in (0, \epsilon) \\ 1 & \|\mathbf{x} - \mathbf{x}^B(t)\|_2 - r^B + \epsilon/2 \leq 0 \end{cases}$$

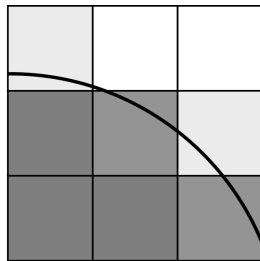
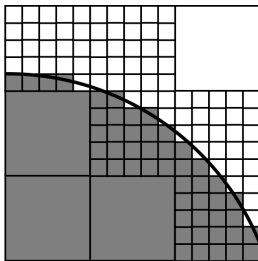
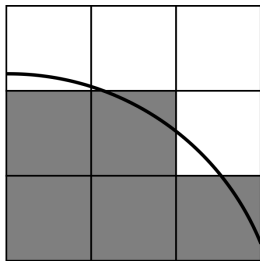


Supersampling

- Every cell, where $\varepsilon(\mathbf{x}, t) \in (0, 1)$, is divided in a subgrid
- Subgrid is constructed with a certain refinement level r
- Binary information of the subgrid is used to estimate $\varepsilon(\mathbf{x}, t)$

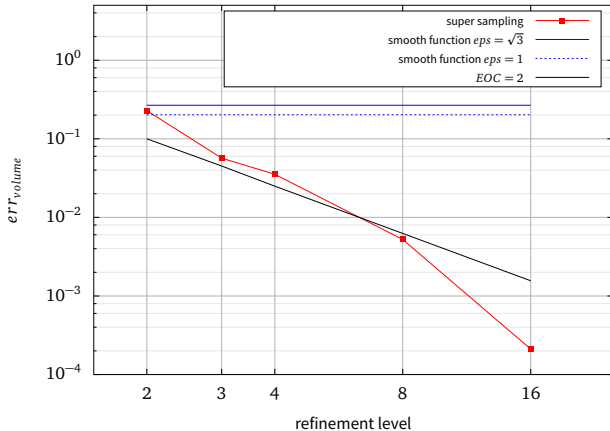
$$\varepsilon(\mathbf{x}, t) = \sum_{i=1}^{r^3} V_{s,i}^{subgrid}(\mathbf{x}, t)$$

$$V_{s,i}^{subgrid}(\mathbf{x}, t) = \begin{cases} 0 & false \\ \frac{1}{r^3} & true \end{cases}$$



Supersampling vs. Smoothed Profile

$$err_{volume} = \frac{1 - \frac{V_r}{V_{analytical}}}{1 - \frac{V_{staircase}}{V_{analytical}}}$$



- Collide and stream

$$f_i(\mathbf{x} + \mathbf{c}_i \Delta t, t + \Delta t) = f_i(\mathbf{x}, t) + B \Omega_i^s(\mathbf{x}, t) + (1 - B) \Omega_i^f(\mathbf{x}, t)$$

- Weighting factor $B(\mathbf{x}, t)$

$$B(\mathbf{x}, t) = \frac{\varepsilon(\mathbf{x}, t) \left(\frac{\tau}{\Delta t} - \frac{1}{2} \right)}{(1 - \varepsilon(\mathbf{x}, t)) + \left(\frac{\tau}{\Delta t} - \frac{1}{2} \right)}$$

- Solid collision operator $\Omega_i^s(\mathbf{x}, t)$

$$\Omega_i^s(\mathbf{x}, t) = \left[f_{\bar{i}}(\mathbf{x}, t) - f_{\bar{i}}^{\text{eq}}(\rho, \mathbf{u}_s) \right] - \left[f_i(\mathbf{x}, t) - f_i^{\text{eq}}(\rho, \mathbf{u}_s) \right]$$

- Fluid collision operator $\Omega_i^f(\mathbf{x}, t)$

$$\Omega_i^s(\mathbf{x}, t) = \frac{\Delta t}{\tau} \left(f_i^{\text{eq}}(\rho_f, \mathbf{U}) - f_i(\mathbf{x}, t) \right)$$

- Total hydrodynamic force $F(t)$

$$F(t) = \frac{(\Delta x)^3}{\Delta t} \sum_{\mathbf{x}_s} \left[B(\mathbf{x}_s, t) \sum_i (\Omega_i^s(\mathbf{x}_s, t) \mathbf{c}_{\bar{i}}) \right]$$

- Boundary torque $T(t)$

$$T(t) = \frac{(\Delta x)^3}{\Delta t} \sum_{\mathbf{x}_s} \left[B(\mathbf{x}_s, t) (\mathbf{x}_s - \mathbf{R}) \times \sum_i (\Omega_i^s(\mathbf{x}_s, t) \mathbf{c}_{\bar{i}}) \right]$$

- Collide and stream

$$f_i(\mathbf{x} + \mathbf{c}_i \Delta t, t + \Delta t) = f_i(\mathbf{x}, t) + \Omega_i^f(\mathbf{x}, t)$$

- Weighting factor $B(\mathbf{x}, t)$

$$B(\mathbf{x}, t) = \varepsilon(\mathbf{x}, t)$$

- Fluid collision operator $\Omega_i^f(\mathbf{x}, t)$

$$\Omega_i^f(\mathbf{x}, t) = \frac{\Delta t}{\tau} (f_i^{\text{eq}}(\rho_f, \bar{\mathbf{u}}) - f_i(\mathbf{x}, t))$$

- Equilibrium velocity $\bar{\mathbf{u}}$

$$\bar{\mathbf{u}} = \mathbf{U} + B(\mathbf{x}, t)(\mathbf{u}_s - \mathbf{U})$$

- Momentum Exchange $g_i(\mathbf{x}, t)$

$$g_i(\mathbf{x}, t) = \mathbf{c}_i f_i(\mathbf{x}, t) + \mathbf{c}_{\bar{i}} f_i(\mathbf{x} + \mathbf{c}_{\bar{i}} \Delta t, t)$$

- Total hydrodynamic force $\mathbf{F}(t)$

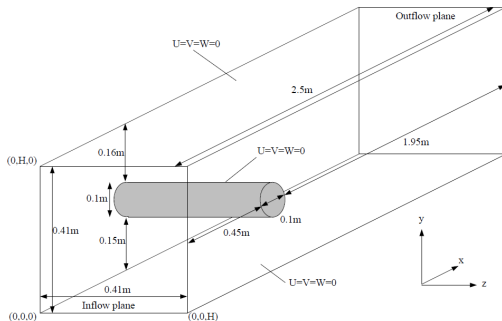
$$\mathbf{F}(t) = \frac{(\Delta x)^3}{\Delta t} \sum_{\mathbf{x}_s} \left[B(\mathbf{x}_s, t) \sum_i g_i(\mathbf{x}, t) \right]$$

- Boundary torque $\mathbf{T}(t)$

$$\mathbf{T}(t) = \frac{(\Delta x)^3}{\Delta t} \sum_{\mathbf{x}_s} \left[B(\mathbf{x}_s, t) (\mathbf{x}_s - \mathbf{R}) \times \sum_i g_i(\mathbf{x}, t) \right]$$

■ Drag and lift coefficients

$$c_D = \frac{2F_w}{\rho \bar{u}^2 D H}, \quad c_L = \frac{2F_a}{\rho \bar{u}^2 D H}$$



C_D	$N = 10$	$N = 20$	$N = 40$
HLBM $\epsilon = 2$	9.83	7.57	6.81
HLBM $\epsilon = 1$	8.15	6.90	6.51
HLBM $\epsilon = 0$	6.90	6.38	6.22
PSM	6.47	6.30	6.24
Bouzidi	6.35	6.22	6.19
Schäfer et al.	6.05-6.25		

$C_L (\times 10^{-2})$	$N = 10$	$N = 20$	$N = 40$
HLBM $\epsilon = 2$	0.00	1.04	1.11
HLBM $\epsilon = 1$	0.00	0.66	0.87
HLBM $\epsilon = 0$	0.00	0.53	0.73
PSM	0.47	0.77	0.84
Bouzidi	13.20	2.50	1.25
Schäfer et al.	0.80-1.00		

- Lattice Boltzmann provides resolved and non resolved approaches to simulate suspensions
- Efficient and accurate volume fraction approximation techniques are challenging
- Various resolved approaches offers a wide range of applications

Thank you for your attention!

